



## PAPER

# Influence of Fe<sub>2</sub>O<sub>3</sub>@reduced graphene oxide nanocomposite on the structural, morphological, and optical features of the polyvinyl alcohol films for optoelectronic applications

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E-mail: [tarek.attia@fsc.bu.edu.eg](mailto:tarek.attia@fsc.bu.edu.eg)**Keywords:** polyvinyl alcohol, reduced graphene oxide, refractive index, linear and nonlinear optical parameters

## Abstract

Polyvinyl alcohol (PVA) films doped with Fe<sub>2</sub>O<sub>3</sub>-reduced graphene oxide (rGO) nanoparticles (NPs) were prepared using a casting-method. Fe<sub>2</sub>O<sub>3</sub> NPs were synthesized via the hydrothermal process, and then Fe<sub>2</sub>O<sub>3</sub> NPs were decorated on the graphene oxide (GO) sheets, where the GO was transferred to rGO during the sonication process. The obtained films were characterized using XRD and FTIR techniques. Optical absorption and transmission data were recorded via a UV–visible spectrophotometer and used to estimate various optical parameters. Increasing the amount of doped Fe<sub>2</sub>O<sub>3</sub>-rGO NPs in PVA decreased the ability of the prepared nanocomposites to allow visible light to pass through them. Solar material protection factor (SMPF) of PVA (11.93%) improved to 99.3%, corresponding to 2 wt% of Fe<sub>2</sub>O<sub>3</sub>-rGO NPs doped in the host matrix. The calculated values of the average refractive index are 2.18, 2.25, 2.36, 2.46, and 2.58, corresponding to 0 wt%, 0.5 wt%, 1 wt%, 1.5 wt%, and 2 wt% of Fe<sub>2</sub>O<sub>3</sub>-rGO NPs doped in PVA. Therefore, the average refractive index showed a dependence on Fe<sub>2</sub>O<sub>3</sub>-rGO NPs. Low energy region ( $E < 2.4$  eV) showed dielectric relaxation time-energy dependent behavior, whereas the high energy region ( $E > 2.4$  eV) displayed dielectric relaxation time- energy independent behavior. Increasing the content of the doped NPs in PVA resulted in lowering the surface and volume energy loss.

## 1. Introduction

Polymer nanocomposite materials have evolved into a vital component of technological advancement and are now an integral part of daily life [1–3]. Nowadays most studies focus on the fabrication of innovative polymer nanocomposite materials with enhanced physical properties. Polymer materials have attracted scientific and industrial interest due to their low cost, biodegradability, easy synthesis, and good optoelectrical properties [4]. These materials are also promising alternative materials for inorganic semiconductors in optoelectronic applications like; photovoltaic cells, Supercapacitors, magneto-optic data storages, gas sensors, and biomedical applications [4–7]. Several works have been done to enhance the performance of the optical properties of the polymer material via incorporating various filler materials with different concentrations and/or different sizes [8]. Optical parameters such as absorbance, reflection, transmission, absorption coefficient, optical band gap, Urbach energy, refractive index, extinction coefficient, optical dielectric constant, optical dielectric loss, and optical conductivity are the key factor to evaluate the material's suitability as a suitable candidate for the optoelectronic applications, which can be enhanced and tuned via the changing of the filler nature, size, and its distribution in a polymer matrix.

Polyvinyl alcohol (PVA) is a biocompatible and thermostable polymer material [9]. It has been used as a host matrix for various filler materials, owing to their good features like; non-corrosive nature, soluble in water, thermal stability, good film forming, and highly transparent film. In our previous work, different nanoparticles (NPs) like  $\text{SiO}_2$  [10],  $\text{ZnO}$  [11],  $\text{Cd}_{0.5}\text{Zn}_{0.5}\text{Fe}_2\text{O}_4$  [12],  $\text{La}_2\text{ZnO}_x$  [13], and  $\text{Mn}_2\text{O}_3$ -rGO [14] have been used as filler materials in PVA polymer matrices to tune and enhance optical properties and to tailor their optical parameters for optical industrial applications.

Badawi *et al* [15] studied the effect of reinforcing different concentrations of  $\text{Fe}_2\text{O}_3$  nanoparticles on the structural and optical properties of the PVA/graphene (G) blend. It was found that the addition of  $\text{Fe}_2\text{O}_3$  NPs to the blend matrix causes a decrease in the optical band gap and an enhancement in the optical dispersion parameters. Also, Badawi *et al* [16] studied the effect of reinforcing 2 wt% from  $\text{Fe}_2\text{O}_3$ ,  $\text{Pb}_2\text{O}_3$ , and  $\text{MnO}_2$  on the structural and optical properties of PVA/rGO blend films. The obtained results revealed that  $\text{Fe}_2\text{O}_3$  has more effect on the decrease in the optical band gap than  $\text{Pb}_2\text{O}_3$ , and  $\text{MnO}_2$ . Aslam *et al* [17] reported an improvement in the optical absorption of the PVA polymer films when reinforced with GO sheets. Recently, researchers have succeeded to decorate various nanoparticles on reduced graphene (rGO) sheets to enhance their performance [18–26]. Previously, our team succeeded in fabricating  $\text{Mn}_2\text{O}_3$ -rGO composite which was used as a filler material in the PVA matrix [14].

As mentioned earlier [15, 16], two blends of polymers, PVA/G and PVA/rGO, have been prepared and then reinforced with different concentrations of  $\text{Fe}_2\text{O}_3$  NPs and different nanoparticle types, respectively. Herein, the  $\text{Fe}_2\text{O}_3$  NPs were firstly synthesized via the hydrothermal process and then the obtained sample was decorated on the graphene oxide (GO) sheets. Moreover, the GO was transferred to rGO during the sonication process which is in good accordance with the literature [19]. Then, the obtained composite,  $\text{Fe}_2\text{O}_3$ -rGO, was reinforced in the host matrix, PVA, with different concentrations. This study intends to investigate the linear and nonlinear optical parameters of PVA films reinforced with different concentrations of  $\text{Fe}_2\text{O}_3$ -rGO nanoparticles. The x-ray diffractometer (XRD), Fourier transfer infrared (FTIR), and UV–visible techniques were used to study the structure and optical properties of the prepared polymer films, respectively. This study provides reporting data about the enhancement of their performance in optoelectronic applications.

## 2. Experimental and methods

### 2.1. Materials

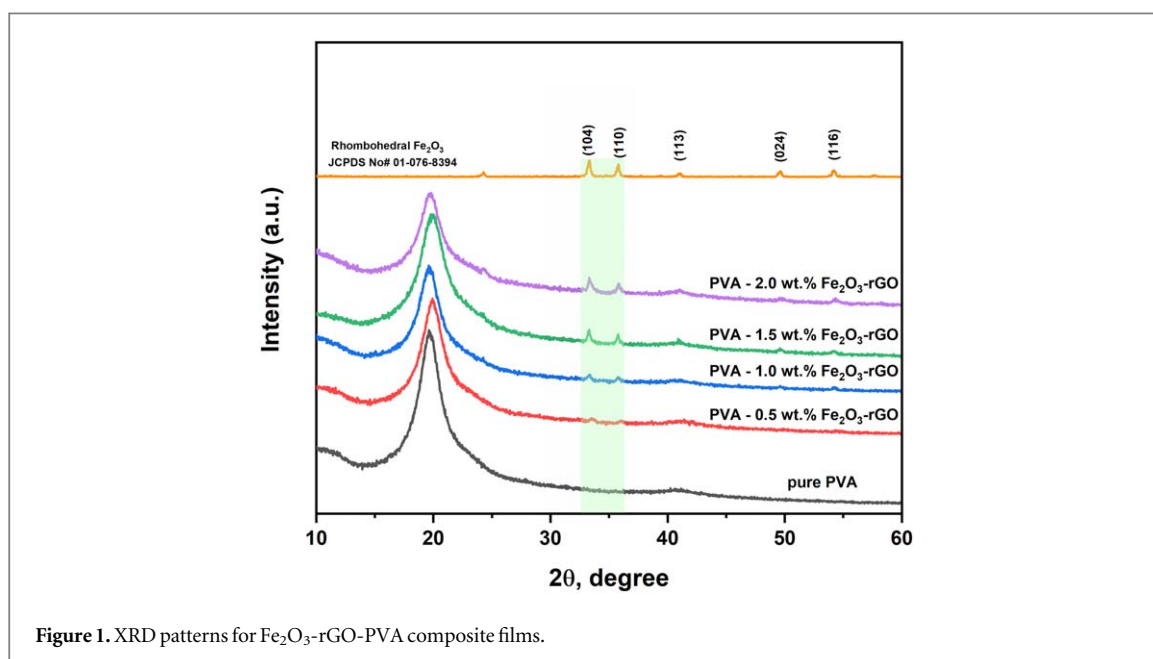
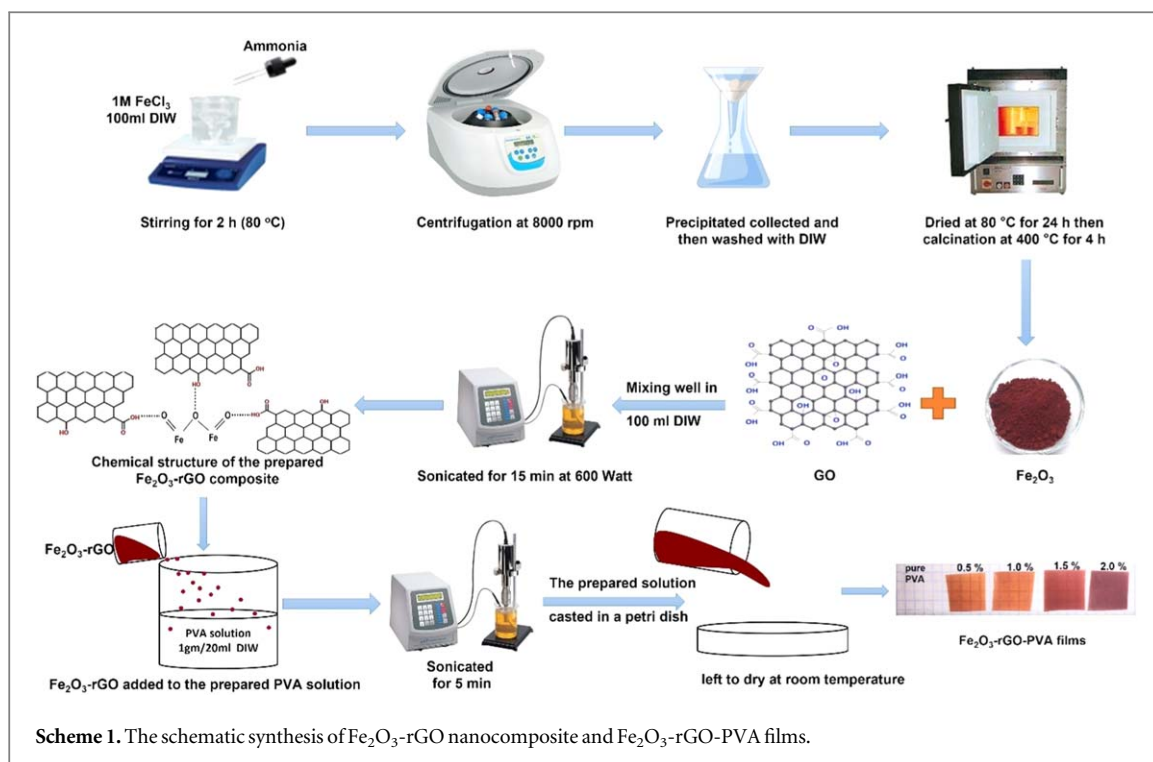
Polyvinyl alcohol (PVA), with a molecular weight  $= 1 \times 10^6$ , was purchased from Qualikem Company.  $\text{Fe}_2\text{O}_3$ -rGO NPs were prepared by the co-precipitation method. 1 M  $\text{FeCl}_3$  in 100 ml deionized water (DIW) was heated to  $80^\circ\text{C}$ , and ammonia was added to the solution with continuous stirring for 2 h. Then, the product was centrifuged (8000 rpm), and the precipitated particles were collected and washed with DIW. The prepared  $\text{Fe}_2\text{O}_3$  NPs were dried for 24 h. (at  $80^\circ\text{C}$ ), followed by calcination for 4 h (at  $400^\circ\text{C}$ ). Finally, the prepared  $\text{Fe}_2\text{O}_3$  NPs were decorated on GO by mixing them in 100 ml DIW and sonicating them for 15 min (GO transferred to rGO during the sonication process [14]).

### 2.2. Preparation of $\text{Fe}_2\text{O}_3$ -rGO-PVA nanocomposite films

Firstly, 5 gm of PVA powder was dissolved in 100 ml distilled water at  $70^\circ\text{C}$ , with the help of magnetic stirrer. Then, different concentrations of  $\text{Fe}_2\text{O}_3$ -rGO NPs were added to the host matrix and distributed via the probe ultrasonic process. The sonication continued for 5 min. Then the solution was kept in a vacuum chamber for an hour to remove any bubbles produced during the sonication process. Finally, the solution was cast in glass petri-dishes and lifted to dry at room temperature and dry atmosphere for a week. The obtained polymer films were kept in a vacuum chamber for a night to remove any residual solvents from the polymer films. The polymer film thickness was measured using a digital micrometer and was found to be about  $100\ \mu\text{m}$ . The chemical structure and schematic synthesis of  $\text{Fe}_2\text{O}_3$ -rGO nanocomposite and  $\text{Fe}_2\text{O}_3$ -rGO-PVA film is shown in scheme 1.

### 2.3. Characterization and setups

The x-ray diffraction pattern of the  $\text{Fe}_2\text{O}_3$ -rGO-PVA films was investigated through Bruker D8 diffractometer,  $\text{Cu-K}\alpha$  radiation  $\lambda = 1.5418\ \text{\AA}$ . Fourier transform infrared (FTIR, Vertex 70-Bruker-Germany) spectrometer was used with range ( $4000$  to  $400\ \text{cm}^{-1}$ ,  $4\ \text{cm}^{-1}$  resolution). The  $\text{Fe}_2\text{O}_3$ -rGO-PVA films morphology was investigated using OLYMPUS-BX51 optical microscope and Carl-Zeiss LEO982 scanning electron microscope (SEM). The UV-visible absorption and transmission spectra were collected via Cary5000-UV spectrophotometer.



### 3. Results and discussions

#### 3.1. XRD analysis

Figure 1 shows the XRD pattern of the pristine sample and the pristine reinforced with different concentrations of  $\text{Fe}_2\text{O}_3$ -rGO NPs. The well-known broad hump at  $19.62^\circ$  is indexed to the (101) diffraction plane of the characteristic peak of semicrystalline PVA film.

With the addition of  $\text{Fe}_2\text{O}_3$ -rGO NPs to the PVA matrix, the hump descends and broadened gradually with increasing the NPs concentration. In addition, the peaks at  $33.3^\circ$  and  $35.8^\circ$ , indexed to the rhombohedral  $\text{Fe}_2\text{O}_3$  NPs (JCPDS No. # 01-076-8394), observed, and became more pronounced at higher  $\text{Fe}_2\text{O}_3$  NPs concentration. This reveals the increase in PVA film amorphousity with the addition of the  $\text{Fe}_2\text{O}_3$ -rGO NPs to the matrix [27, 28]. In addition, the  $\text{Fe}_2\text{O}_3$  peaks at higher concentration became more pronounced and vanishes of it at lower concentration is an indication of the increase of the particle size because of the agglomeration with rising

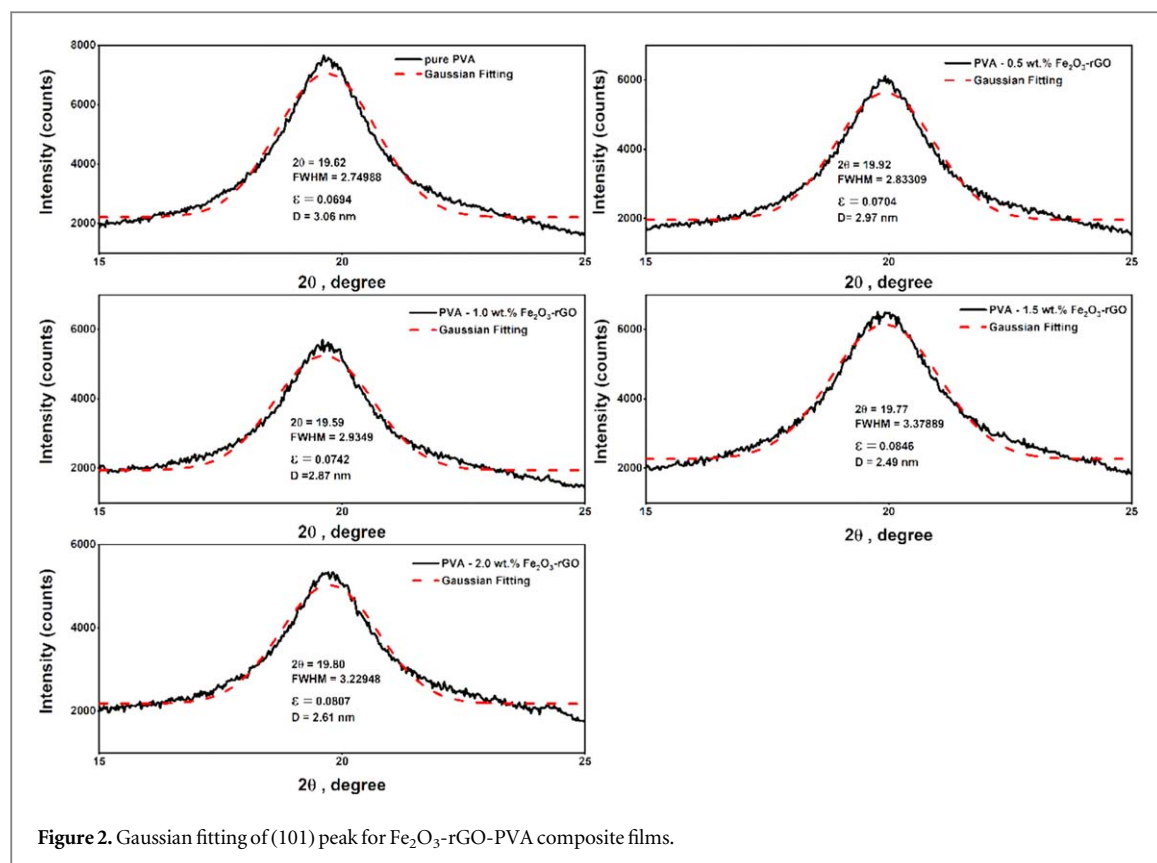


Figure 2. Gaussian fitting of (101) peak for  $\text{Fe}_2\text{O}_3$ -rGO-PVA composite films.

Table 1. Geometrical parameters of  $\text{Fe}_2\text{O}_3$ -rGO-PVA composite films.

Sample	FWHM (degree)	$2\theta$ , degree	D (nm)	$\epsilon$	$\delta (\text{nm}^{-2})$
Pure PVA	2.74	19.62	3.06	0.069	0.107
0.5%	2.83	19.92	2.97	0.070	0.113
1.0%	2.93	19.59	2.87	0.074	0.121
1.5%	3.38	19.77	2.49	0.085	0.161
2.0%	3.23	19.80	2.61	0.081	0.147

of  $\text{Fe}_2\text{O}_3$ -rGO NPs concentration. Furthermore, the H-bonds formed between the  $\text{Fe}_2\text{O}_3$ -rGO NPs and the hydroxyl groups of PVA molecules. Recently, Badawi *et al* [15] concluded that the addition of  $\text{Fe}_2\text{O}_3$  to the PVA/graphene decreases the PVA semi-crystallinity.

Gaussian fitting of the characteristic peak (101) of PVA films was used to estimate the crystallite size (D), the internal strain ( $\epsilon$ ), and dislocation density ( $\delta$ ), as shown in figure 2. This can be calculated by the following equations [14],

$$D = \frac{0.9 \lambda}{\beta \cos \theta} \quad (1)$$

$$\epsilon = \frac{\beta}{4 \tan \theta} \quad (2)$$

$$\delta = \frac{1}{D^2} \quad (3)$$

where  $\beta$  is the full width at half maximum (FWHM) and  $\theta$  is the Bragg's angle. The obtained data are presented in table 1 and shown in figure 2.

Obviously, as the  $\text{Fe}_2\text{O}_3$ -rGO NPs concentration in the PVA matrix increases the D value decreases. This is caused due to the incorporation of the  $\text{Fe}_2\text{O}_3$ -rGO NPs into the PVA matrix causing a decrement in the hydrogen bonding in the PVA matrix and forming of hydrogen bonding between the  $\text{Fe}_2\text{O}_3$ -rGO NPs and the PVA molecules [29–31]. In contrary, the  $\epsilon$  and  $\delta$  values are found to increase with the increase in the  $\text{Fe}_2\text{O}_3$ -rGO concentration. The high values of  $\epsilon$  and  $\delta$  confirms the decrease in polymer films crystallinity [27, 32].

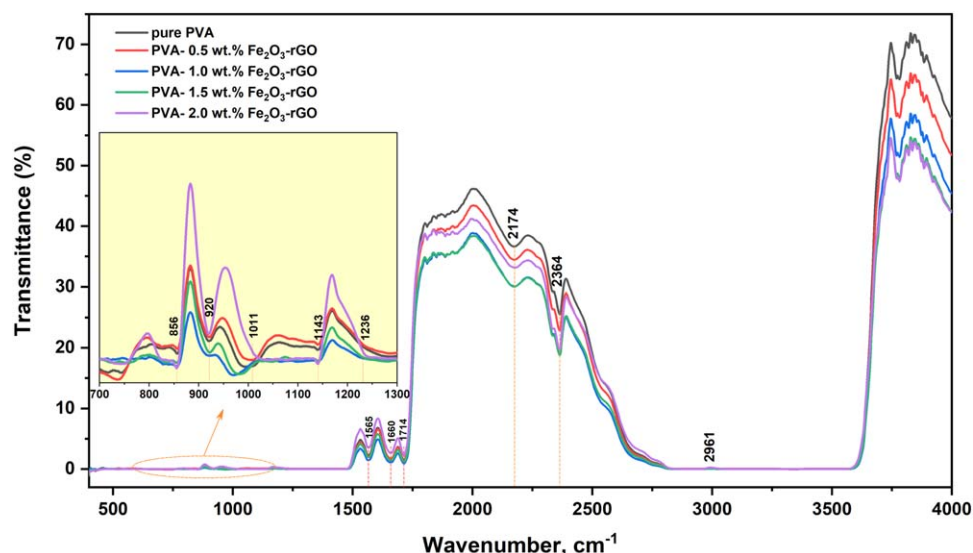


Figure 3. FTIR spectra for  $\text{Fe}_2\text{O}_3$ -rGO-PVA composite films.

### 3.2. FTIR and optical microscope analysis

Figure 3 shows the FTIR spectra of  $\text{Fe}_2\text{O}_3$ -rGO-PVA films. It shows a broad band at  $3000\text{--}3600\text{ cm}^{-1}$  which is assigned to O–H stretching vibration mode of PVA molecules. The absorption band at  $2961\text{ cm}^{-1}$  is related to C–H asymmetric stretching mode. Furthermore, the bands  $1714$ ,  $1660$ ,  $1565\text{ cm}^{-1}$  are representing the C=O vibrational stretching, C=C stretching vibration, and (O–H&C–H) bending modes, respectively [33–35].  $1143\text{ cm}^{-1}$  related to the C–O–C vibrational stretching mode of acetyl group in PVA matrix, and  $1236\text{ cm}^{-1}$  related to  $-\text{CH}_2$  wagging. While the bands at  $856$ ,  $1011\text{ cm}^{-1}$  are related to the C–H rocking vibrations and C–O stretching of acetyl groups, respectively. The band at  $920\text{ cm}^{-1}$  is related to the epoxides in GO [31]. The bands intensity was changed with the addition of  $\text{Fe}_2\text{O}_3$ -rGO NPs in the PVA matrix which may be caused by the formation of the hydrogen bonds between the  $\text{Fe}_2\text{O}_3$ -rGO NPs and the PVA molecules [13]. In addition, the incorporation of  $\text{Fe}_2\text{O}_3$ -GO NPs in the PVA matrix causes a decrement in the hydrogen bonds between the PVA molecules and produces hydrogen bonds between  $\text{Fe}_2\text{O}_3$ -GO NPs and PVA molecules [31]. The band position was slightly changed as in the band at  $856\text{ cm}^{-1}$  for pure PVA shifted to lower value at  $854\text{ cm}^{-1}$  for 2 wt%  $\text{Fe}_2\text{O}_3$ -rGO-PVA film. This slight shift is an indication of the formation of hydrogen bonds [36].

The morphological study of the polymer films surface was investigated using the optical microscope. In figure 4 a smooth surface was observed in the bright field image of the pure PVA. With the addition of  $\text{Fe}_2\text{O}_3$ -rGO NPs in the PVA matrix, particles, small spherical, and big irregular shapes were seen in the polymer films. The big particles represent clusters of  $\text{Fe}_2\text{O}_3$ -rGO NPs, which formed due to aggregations caused by the magnetic inductions between the  $\text{Fe}_2\text{O}_3$ -rGO NPs each other. This aggregation increases with increasing the  $\text{Fe}_2\text{O}_3$ -rGO NPs concentrations in the PVA matrix. Figures 4(f)–(g) shows the surface morphology of PVA films doped with  $\text{Fe}_2\text{O}_3$ -rGO NPs. PVA films look rough, and the white spots represent the embedded  $\text{Fe}_2\text{O}_3$ -rGO NPs. In addition, these NPs grow in size and are denser at 2.0 wt% (figure 4(g)) as compared with 0.5 wt% (figure 4(f)). It is clear from the histogram shown in figure 4(h) that the minimum and maximum particle sizes are about 44 and 413 nm for PVA-0.5 wt%  $\text{Fe}_2\text{O}_3$ -rGO film and about 60 and 456 nm for PVA-2.0 wt%  $\text{Fe}_2\text{O}_3$ -rGO film, respectively. The growth in the crystallite size was caused due to the aggregations, which are in good accordance with the XRD data.

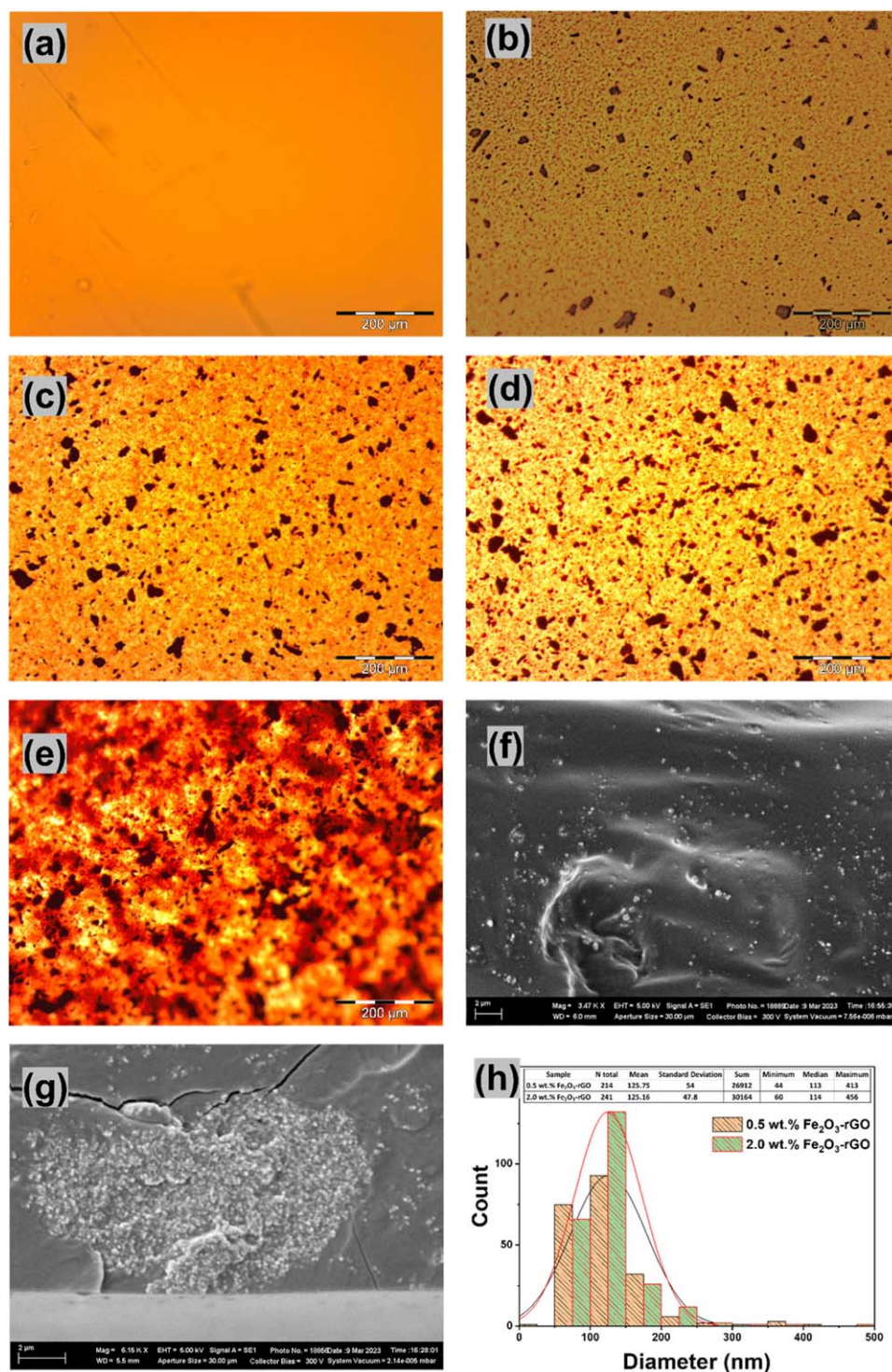
### 3.3. Optical properties

#### 3.3.1. Skin depth, interband transitions, and urbach energy

UV-vis spectrometer was used to measure the transmission and absorption of the nanocomposites of  $\text{Fe}_2\text{O}_3$ -rGO-PVA in a range of wavelengths from 200 nm to 1000 nm, as seen in figures 5(a), (b), respectively. The host matrix (PVA) has a transmission value of 89% at 400 nm, which decreased to 11%, 3%, 0.9% and 0.4%, corresponding to the added amount of  $\text{Fe}_2\text{O}_3$ -rGO NPs of 0.5 wt%, 1 wt%, 1.5 wt%, and 2 wt%, respectively. Therefore, increasing the doped content of  $\text{Fe}_2\text{O}_3$ -rGO NPs in the host matrix caused an obvious decrease in the transmission due to acting these NPs as scattering centers, in satisfactory accordance with the relevant literature [37].

Plotting absorption of nanocomposite divided by the absorption of PVA versus wavelength, as seen in figure 5(c) illustrated an obvious absorption peak in the visible region ( $\lambda > 400\text{ nm}$ ), which was not clear in the

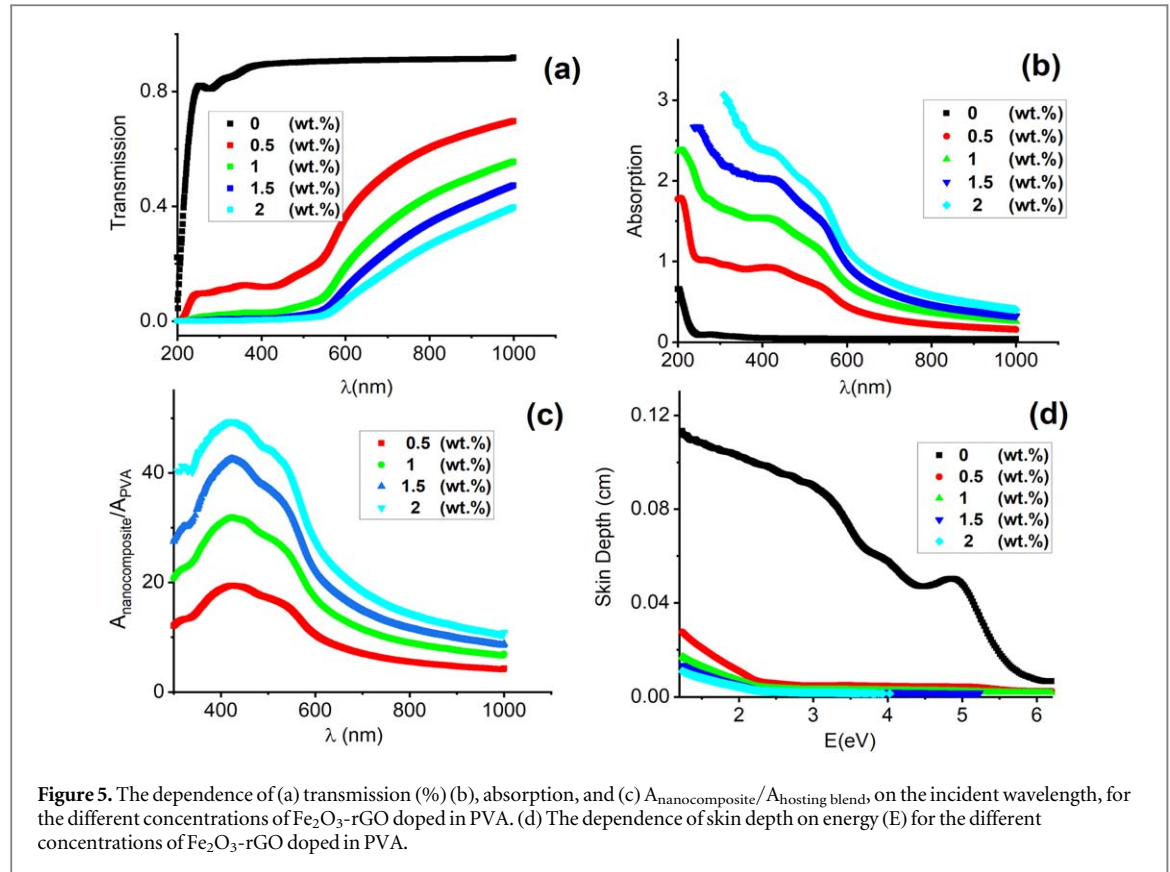




**Figure 4.** Optical microscope images (a-e), SEM images (f-g), and the histogram (h) of PVA- Fe<sub>2</sub>O<sub>3</sub>-rGO films. (a) 0, (b) & (f) 0.5, (c) 1.0, (d) 1.5, and (e) & (g) 2.0 wt% Fe<sub>2</sub>O<sub>3</sub>-rGO.

absorption spectra of each nanocomposite, see figure 5(b). The positions of the observed peaks in figure 5(c) are at wavelengths of 425 nm, 423 nm, 421 nm, and 416 nm corresponding to the doped amounts of Fe<sub>2</sub>O<sub>3</sub>-rGO of 0.5 wt%, 1 wt%, 1.5 wt%, and 2 wt%, respectively. The clear peaks in figure 5(c) correspond to Fe<sub>2</sub>O<sub>3</sub> NPs, in good agreement with a published work [38].

The distance that an electromagnetic wave is able to go into a polymeric substance is referred to as the skin depth ( $\delta$ ) [39]. The value of the skin depth can be calculated by taking the reciprocal of the absorption coefficient ( $\alpha$ ), which represents the ratio of the absorption to the sample thickness). For the different nanocomposites of Fe<sub>2</sub>O<sub>3</sub>-rGO-PVA, the dependence of the skin depth ( $\delta$ ) on wavelength is depicted in figure 5(d). Increasing the content of the Fe<sub>2</sub>O<sub>3</sub>-rGO NPs in the host matrix of PVA resulted in a decrease in the skin depth of the



**Figure 5.** The dependence of (a) transmission (%) (b), absorption, and (c)  $A_{\text{nanocomposite}}/A_{\text{hosting blend}}$ , on the incident wavelength, for the different concentrations of  $\text{Fe}_2\text{O}_3$ -rGO doped in PVA. (d) The dependence of skin depth on energy ( $E$ ) for the different concentrations of  $\text{Fe}_2\text{O}_3$ -rGO doped in PVA.

nanocomposites. This behavior can be attributed to increasing the absorption and decreasing the transparency by increasing the NPs content.

The absorption spectrum fitting (ASF) method was successfully used to estimate the optical band gap based on the absorption spectra according to the following relation [40, 41]:

$$\alpha(\lambda) = Z(hc)^{-1+m} \lambda(\lambda^{-1} - \lambda_g^{-1})^m \quad (4)$$

Where  $\alpha$  is the absorption coefficient,  $h$  is Plank's constant,  $Z$  is a constant,  $c$  is the light speed, and  $\lambda_g$  represents the wavelength where the optical gap occurs. The index ' $m$ ' defines the kind of electronic transitions responsible for the absorption process. The values of  $m$  can be  $1/3$  and  $3$  correspond to direct and indirect forbidden transitions, respectively, whereas the values of  $0.5$  and  $2$  are related to the direct and indirect allowed transitions, respectively. The best fit was achieved for  $m = 0.5$ , implying that direct allowed transition was demonstrated for all samples. The value of  $\lambda_g^{-1}$  shown in figure 6(a) was obtained by extrapolating the linear portion of the plot of  $(A\lambda^{-1})^{1/m}$  against  $\lambda^{-1}$  at the zero value of  $(A\lambda^{-1})^{1/m}$ . The values of ASF optical gap ( $E_{\text{gab}}^{\text{ASF}}$ ) were calculated by the product of  $1239.83$  times  $\lambda_g^{-1}$  and listed in table 2 for all of the samples under investigation.

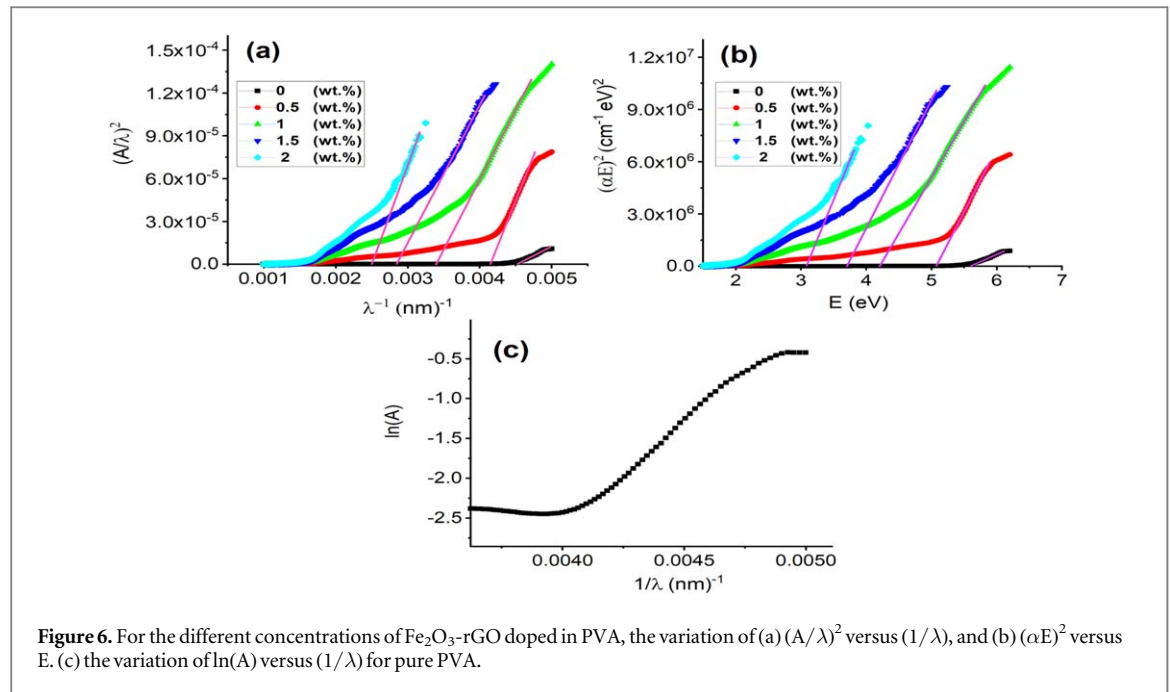
The optical gap energies were also calculated using Tauc's formula, which was derived from the dependence of the absorption coefficient on the photon energy as follows [38, 39]:

$$\alpha E = B(E - E_{\text{gab}}^{\text{Tauc}})^m \quad (5)$$

where  $B$  is a constant that depends on the transition probability,  $E_{\text{gab}}^{\text{Tauc}}$  is the Tauc band gap energy,  $E$  is the photon energy (eV), and  $m$  expresses the nature of the electronic transition. The dependence of  $(\alpha E)^{1/m}$  on the photon energy for the different nanocomposites is illustrated in figure 6(b). The intercept on the energy axis of the best linear fit of the data in figure 6(b), for  $m = 0.5$ , was used to derive the values of  $E_{\text{gab}}^{\text{Tauc}}$ . Table 2 reported the values of  $E_{\text{gab}}^{\text{Tauc}}$  for all the investigated samples which were obtained from the direct allowed transition.

The values of  $E_{\text{gab}}^{\text{Tauc}}$  are  $5.956$  eV,  $5.061$  eV,  $4.298$  eV,  $3.694$  eV, and  $3.076$  eV, whereas  $E_{\text{gab}}^{\text{ASF}}$  has values of  $5.579$  eV,  $5.083$  eV,  $4.091$  eV,  $3.472$  eV, and  $3.099$  eV, corresponding to  $0$  wt%,  $0.5$  wt%,  $1$  wt%,  $1.5$  wt%, and  $2$  wt% of  $\text{Fe}_2\text{O}_3$ -rGO NPs doped in PVA. There is a similarity between the two methodologies used to calculate the optical gap energy of nanocomposites of  $\text{Fe}_2\text{O}_3$ -rGO-PVA.

The energy gap of bulk  $\text{Fe}_2\text{O}_3$  ( $2.1$  eV) [42] is less than the energy gap of the prepared  $\text{Fe}_2\text{O}_3$  NPs only ( $3.67$  eV) or  $\text{Fe}_2\text{O}_3$ -rGO NPs only ( $4.08$  eV), as deduced in previous work [43]. This result is due to the influence of quantum confinement in agreement with the literature [43]. In addition, increasing the content of the doped  $\text{Fe}_2\text{O}_3$ -rGO NPs in PVA resulted in decreasing the optical gap energy of the host matrix. This result can be due to



the increase in the size of Fe<sub>2</sub>O<sub>3</sub>-rGO NPs due to aggregation and the modification in the structure of PVA by the added NPs [44, 45]. This finding is in good agreement with XRD and SEM results. Additionally, as stated in the literature [37], reducing the gap energy by adding more NPs indicates that the Fe<sub>2</sub>O<sub>3</sub>-rGO-PVA nanocomposites are becoming more semiconducting in nature.

Urbach's law [46], which relates the absorption coefficient to the incident photon energy, was successfully revised in the ASF model, leading to the following modified formula [47]:

$$A(\lambda) = P \exp\left(\frac{hc\lambda^{-1}}{E_{Tail}}\right) \quad (6)$$

Where  $P$  is a constant, and  $E_{Tail}$  is the width of the tail of localized states (Urbach energy), when an optical transition occurs between a localized tail state near the valence band and an extended conduction band state above the mobility edge. The inverse of the slope, generated by the best linear fit of the  $\ln A - \lambda^{-1}$  plot, was multiplied with 1239.83 to determine  $E_{Tail}$ . The dependence of  $\ln A$  on  $\lambda^{-1}$  is shown in figure 6(c) for PVA, as an example for all the examined samples. As shown in table 2, increasing the content of Fe<sub>2</sub>O<sub>3</sub>-rGO NPs doped in PVA resulted in increasing  $E_{Tail}$  of PVA from 0.465 eV to be 0.820 eV for the highest content (2 wt%). Consequently, the defects increased in the nanocomposites by increasing the content of the doped NPs. As displayed in table 2, the values of  $E_{Tail}$ , for all the studied samples, illustrate an opposite behavior of the optical gaps [ $(E_{gab}^{ASF})$  and  $(E_{gab}^{Tauc})$ ]. Table 2 contains a comparison between the optical gap energy and Urbach energy acquired from various earlier publications [37–39, 48–50] and the results of the current investigation.

The strength of electron-phonon ( $E_{e-ph}$ ) interactions is termed by the steepness parameter ( $S$ ) ( $E_{e-ph} = 2/3S$  [51]), which causes a broadening of the absorbance edge:

$$S = \frac{T_{room} k_B}{E_{Tail}} \quad (7)$$

where  $k_B$  is the Boltzmann constant, and  $T_{room}$  is the room temperature in Kelvin.  $E_{e-ph}$  rose from 11.90 for PVA to 20.94 for 2 wt% of the doped Fe<sub>2</sub>O<sub>3</sub>-rGO NPs in PVA, as shown in table 2. This behavior can be a result of increasing the filled bands by increasing the added amounts of the doped Fe<sub>2</sub>O<sub>3</sub>-rGO NPs [52].

Obtaining the number of carbon atoms per molecule based on the values of the calculated energy of the optical gap ( $E_{gab}^{Tauc}$ ), was proposed by Fink *et al* [53] using the following relation:

$$E_{gap}^{Tauc} = \frac{34.3}{\sqrt{N}} \quad (8)$$

where  $N$  is the number of carbon atoms in the carbonaceous cluster. Table 2 shows that the values of  $N$  are 38, 46, 64, 86, and 124, corresponding to 0 wt%, 0.5 wt%, 1 wt%, 1.5 wt%, and 2 wt%. of Fe<sub>2</sub>O<sub>3</sub>-rGO NPs doped in PVA. Therefore, increasing the added amount of the doped Fe<sub>2</sub>O<sub>3</sub>-rGO NPs can increase the conjugation in the monomer unit of the host matrix. This result agrees well with a previously published work [54].



**Table 2.** Tauc optical gap ( $E_{\text{gab}}^{\text{Tauc}}$ ), ASF optical gap ( $E_{\text{gab}}^{\text{ASF}}$ ), tailing state energy ( $E_{\text{Tail}}$ ), the steepness parameter (S), the strength of electron-phonon interactions ( $E_{\text{e-ph}}$ ), carbon atoms per conjugated length (N), solar material protection factor (SMPF), and visible solar transmittance ( $T_{\text{vis}}$ ).

Optical Parameter	Current Work					Previously published works				
	Content of Fe <sub>2</sub> O <sub>3</sub> -rGO (wt%) doped in PVA					PVA [48]	0.5 & 1wt% of Fe <sub>2</sub> O <sub>3</sub> doped in PVA/PEG [37]	1 wt% & 5 wt% of Fe <sub>2</sub> O <sub>3</sub> NPs doped in chitosan [38]	0% & 0.012% (w/w) of rGO doped in PVA [49]	0% & 2% of rGO doped in PVA [50]
		0	0.5	1	1.5	2				
$E_{\text{direct}}^{\text{ASF}}$ (eV)	5.579	5.083	4.091	3.472	3.099					
$E_{\text{direct}}^{\text{Tauc}}$ (eV)	5.595	5.061	4.298	3.694	3.076	5.31	5.15 & 4.83	3.14–2.80	5.29 & 3.59	3.57 & 2.72
$E_{\text{Tail}}$ (eV)	0.465	0.740	0.762	0.809	0.820	0.20	1.10 & 1.95		0.20 & 0.85	
$S \times 10^{-2}$	5.599	3.524	3.426	3.223	3.183					
$E_{\text{e-ph}}$	11.90	18.92	19.42	20.68	20.94					
N	38	46	64	86	124					
$T_{\text{vis}}$ (%)	90.4	25.9	11.4	6.15	3.64					
SMPF (%)	11.3	85.9	95.9	98.5	99.3					

### 3.3.2. Solar protection factors

#### 3.3.2.1. Visible solar transmittance ( $T_{vis}$ )

The visible solar transmittance ( $T_{vis}$ ) gives an indication of the solar energy transmittance. As inferred by the following equation, calculating  $T_{vis}$  can provide valuable information regarding the skin's degree of protection:

$$T_{vis} = 1 - \frac{\sum_{\lambda=380 \text{ nm}}^{\lambda=780 \text{ nm}} \Delta \lambda T_{\lambda} D_{\lambda} V(\lambda)}{\sum_{\lambda=380 \text{ nm}}^{\lambda=780 \text{ nm}} \Delta \lambda D_{\lambda} V(\lambda)} \quad (9)$$

where  $D_{\lambda}$  is the relative spectral distribution of illuminant D65 [55, 56],  $T(\lambda)$  is the spectral transmittance of the studied nanocomposites obtained from figure 5(a),  $\lambda$  is the wavelength,  $\Delta \lambda$  is the wavelength interval ( $\Delta \lambda = 10$  in a range of wavelengths from 380 nm to 780 nm), and  $\Delta \lambda D_{\lambda} V(\lambda)$  are extracted from published literature [57], and  $V(\lambda)$  is the luminous spectral efficiency for photonic vision provided by the standard observer for photometry [56, 58]. The high transmission of visible light through the studied samples was detected in the host matrix of PVA, which has the highest  $T_{vis}$  (90.4%), as reported in table 2. Table 2 shows that increasing the amount of doped  $\text{Fe}_2\text{O}_3$ -rGO NPs in PVA decreased the ability of the prepared nanocomposites to allow visible light to pass through them. The value of  $T_{vis}$  given by 2 wt% of  $\text{Fe}_2\text{O}_3$ -rGO NPs doped in PVA was the lowest possible value (3.64%).

#### 3.3.2.2. Solar material protection factor (SMPF)

The solar material protection factor (SMPF) is a measure of the resistance of PVA and its nanocomposites against deterioration caused by Sunlight. SMPF was calculated from the following formula [57]:

$$SMPF = 1 - \frac{\sum_{\lambda=300 \text{ nm}}^{\lambda=600 \text{ nm}} C_{\lambda} S_{\lambda} \Delta \lambda T(\lambda)}{\sum_{\lambda=300 \text{ nm}}^{\lambda=600 \text{ nm}} C_{\lambda} S_{\lambda} \Delta \lambda} \quad (10)$$

Where  $T(\lambda)$  is the spectral transmittance of the studied samples obtained from figure 5(a),  $S_{\lambda}$  is the relative spectral distribution of Sunlight [56, 59],  $C_{\lambda} = \exp(-0.012\lambda)$ ,  $C_{\lambda} S_{\lambda} \Delta \lambda$  are obtained from a published reference [57], and  $\Delta \lambda$  is the wavelength interval ( $\Delta \lambda = 5$  in a range of 300 nm – 400 nm and  $\Delta \lambda = 10$  in a range of 400 nm–600 nm). The tabulated values of SMPF, for all the samples under study in table 2, show that SMPF of PVA (11.93%) improved to 99.3%, corresponding to 2 wt% of  $\text{Fe}_2\text{O}_3$ -rGO NPs doped in the host matrix.

#### 3.3.3. Refractive Index ( $n$ ) and attenuation coefficient ( $k$ )

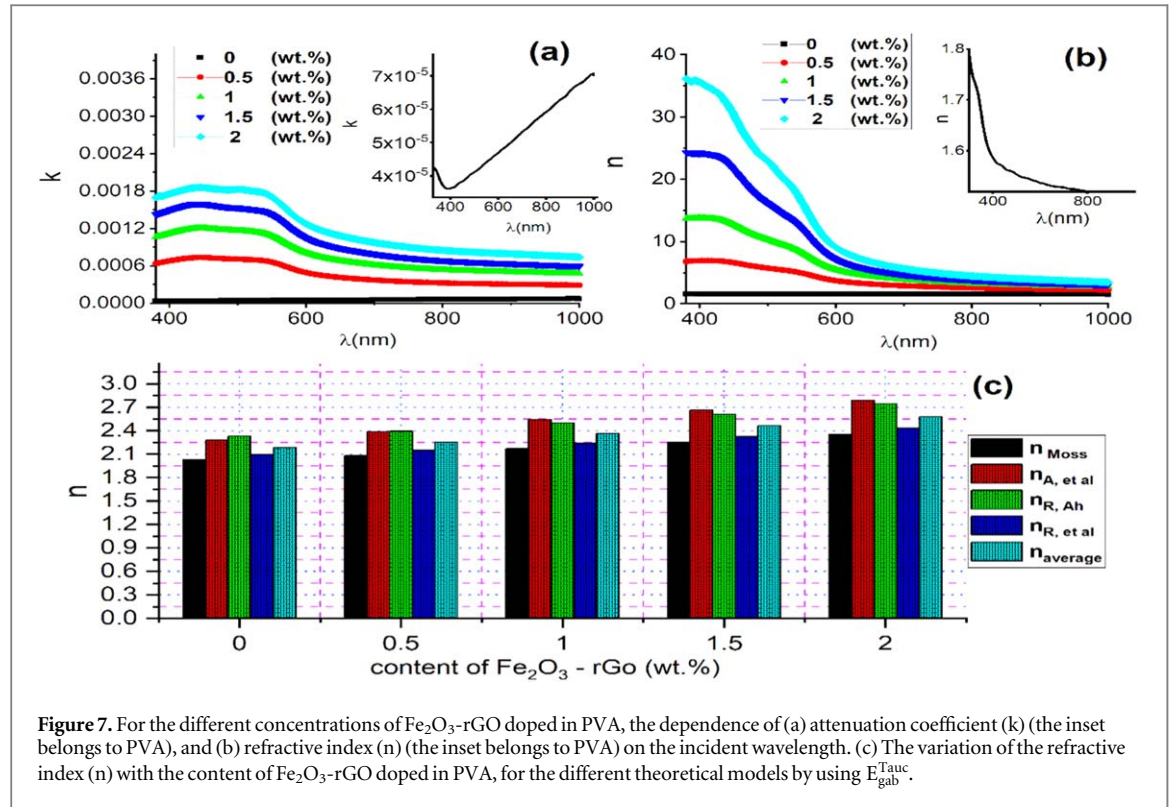
One of the most interesting things about polymeric materials is that their refractive index has a strong relationship with the local electric field and the ability of the electrons inside the material to be polarized. Additionally, the dispersion behavior of the refractive index has an impact on the design of optical communication systems as well as a wide variety of optical instruments. The refractive index was calculated using Fresnel's law as follows [44, 60]:

$$n = \left( \frac{1 + R}{1 - R} \right) + \sqrt{\frac{4R}{(1 - R)^2} - k^2} \quad (11)$$

Where  $R$  is the reflection which was calculated from the relation of  $R = (1 - \text{Transmission} \times \exp(\text{Absorption})^{0.5})$  [5], and  $k$  is the extinction coefficient that depends on the absorption coefficient ( $\alpha$ ) and the incident wavelength ( $\lambda$ ), according to the relation of  $k = \frac{\alpha \lambda}{4\pi}$  [61]. The dependence of the extinction coefficient ( $k$ ) on wavelength for all samples is plotted in figure 7(a). Increasing the content of  $\text{Fe}_2\text{O}_3$ -rGO NPs in PVA enhanced the values of the extinction coefficient ( $k$ ). In addition, increasing the incident wavelength caused an increase in the values of the extinction coefficient ( $k$ ) of the host matrix (PVA) as shown in the inset of figure 7(a), which is attributed to the interaction between the light and the material of the host matrix. This behavior disappeared upon the addition of the doped  $\text{Fe}_2\text{O}_3$ -rGO NPs in agreement with some published studies [37, 48].

Dispersion curves are observed in figure 7(b), for all the content of  $\text{Fe}_2\text{O}_3$ -rGO NPs in PVA, resulting from the slow decrease of the refractive index ( $n$ ) with increasing the incident wavelength. However, the refractive index of PVA suffered from a sharp decrease, as seen in the inset of figure 7(b), typically as reported in the literature for PVA [39]. Increasing the refractive index of PVA by increasing the content of the doped  $\text{Fe}_2\text{O}_3$ -rGO NPs in PVA can be attributed to increasing the number of particles in the material of the nanocomposites. Therefore, the interaction between the incident light and a large amount of particles will increase the refraction, and as a result, the nanocomposites of  $\text{Fe}_2\text{O}_3$ -rGO-PVA will have a higher refractivity [62].

Understanding the relationship between the optical bandgap energy ( $E_g$ ) and the linear refractive index can help researchers understand the band structure of samples and determine more optical parameters. In this study, several models were used to explore the relationship between the energy gap and refractive index, including Moss ( $n_M$ ) [63], Anani, *et al* ( $n_A$ , *et al*) [64], Reddy-Ahammed ( $n_R$ ,  $A_h$ ) [65], Ravindra, *et al* ( $n_R$ , *et al*) [66],



and average refractive index ( $n_{\text{average}}$ ):

$$n_M^4 = \frac{95 \text{ (eV)}}{E_g} \quad (12)$$

$$n_{A, et al} = 3.4 - 0.2 E_g \quad (13)$$

$$n_{R, Ah}^4 = 154(-0.365 + E_g)^{-1} \quad (14)$$

$$n_{R, et al}^4 = \frac{108 \text{ (eV)}}{E_g} \quad (15)$$

$$n_{\text{average}} = [n_M + n_{A, et al} + n_{R, Ah} + n_{R, et al}] \div 4 \quad (16)$$

By using the estimated values of  $E_{\text{direct}}^{\text{Tauc}}$ , and these different models (equation (8)–(12)), the values of the linear refractive index ( $n$ ) were calculated for all the nanocomposites under study, as displayed in figure 7(c). The various models displayed approximations of the calculated linear refractive index for each sample. The calculated values of the average refractive index ( $n_{\text{average}}$ ) are 2.18, 2.25, 2.36, 2.46, and 2.58, corresponding to 0 wt%, 0.5 wt%, 1 wt%, 1.5 wt%, and 2 wt% of  $\text{Fe}_2\text{O}_3$ -rGO NPs doped in PVA. Therefore, the average refractive index ( $n_{\text{average}}$ ) showed a dependence on  $\text{Fe}_2\text{O}_3$ -rGO NPs as it increased by increasing the amount of the doped NPs in PVA.

### 3.3.4. Optical conductivity ( $\sigma_{\text{opt}}$ )

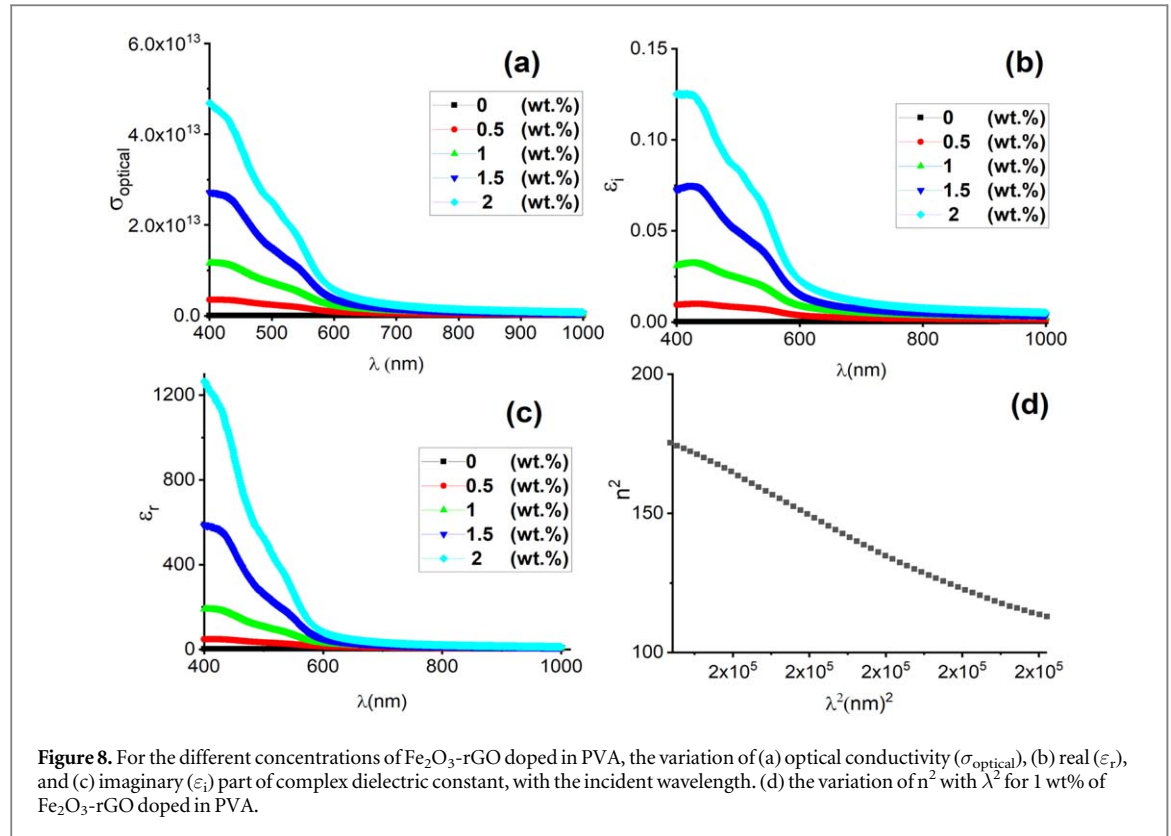
Determining a material's optical conductivity ( $\sigma_{\text{opt}}$ ) can help one comprehend the electronic states that the material possesses, according to the following relation [67]:

$$\sigma_{\text{optical}} = \frac{cn\alpha}{4\pi} \quad (17)$$

where  $c$  is the light velocity,  $\alpha$  is the absorption coefficient,  $n$  is the refractive index, and  $\lambda$  is the wavelength. An obvious development in the value of optical conductivity ( $\sigma_{\text{opt}}$ ) of PVA was observed in figure 8(a) by increasing the doped amount of  $\text{Fe}_2\text{O}_3$ -rGO NPs in PVA. Figure 8(a) shows that at  $\lambda > 750$  nm, the optical conductivity ( $\sigma_{\text{opt}}$ ) is wavelength independent. However, increasing the absorption in the region of wavelength less than 750 nm, as seen before in figure 5(b), caused an improvement in the values of  $\sigma_{\text{opt}}$  for all studied samples.

### 3.3.5. Optical dielectric constants

The optical dielectric constant is directly related to the density of states within the band; therefore, it is beneficial for understanding the band structure of tested samples. The summation of the real part ( $\epsilon_r$ ) and the imaginary



**Figure 8.** For the different concentrations of Fe<sub>2</sub>O<sub>3</sub>-rGO doped in PVA, the variation of (a) optical conductivity ( $\sigma_{\text{optical}}$ ), (b) real ( $\epsilon_r$ ), and (c) imaginary ( $\epsilon_i$ ) part of complex dielectric constant, with the incident wavelength. (d) the variation of  $n^2$  with  $\lambda^2$  for 1 wt% of Fe<sub>2</sub>O<sub>3</sub>-rGO doped in PVA.

part ( $\epsilon_i$ ) produces the complex dielectric function ( $\epsilon^*$ ) where  $\epsilon^* = \epsilon_r + i\epsilon_i$  [68]. The dissipative rate of electromagnetic wave propagation in the medium is due to the imaginary part. However, the real portion of the dielectric constant is proportional to the dispersion, which depends primarily on the motion of electrons in the optical medium as light travels. The real and imaginary parts of the complex dielectric function rely on the refractive index ( $n$ ), and extinction coefficient ( $k$ ) according to the following relations [44, 48]:

$$\epsilon_r = n^2 - k^2, \quad \epsilon_i = 2nk \quad (18)$$

For all the samples being examined, figures 8(b) and (c) illustrate the dependence of  $\epsilon_i$  and  $\epsilon_r$  on the incident wavelength, respectively. Interactions between the incident photon and free electrons in the different nanocomposites can explain the wavelength dependence of the real and imaginary parts of the dielectric constant at  $\lambda < 750$  nm [69]. On the other hand, the values of  $\epsilon_r$  become wavelength independent at higher wavelengths ( $\lambda > 750$  nm) due to the incapability of the molecules of the samples to follow the fluctuations of the incident field [70]. It can be understood from the behavior of  $\epsilon_i$  that the optical loss becomes higher at low wavelength compared to that at high wavelength. As the dielectric constants are dependent on the attenuation coefficient and refractive index, increasing the doped Fe<sub>2</sub>O<sub>3</sub>-rGO NPs in PVA increased the dielectric constants, including  $\epsilon_i$  and  $\epsilon_r$ . This behavior was identical to that observed in figures 7(a) and (b) for the attenuation coefficient and refractive index, respectively.

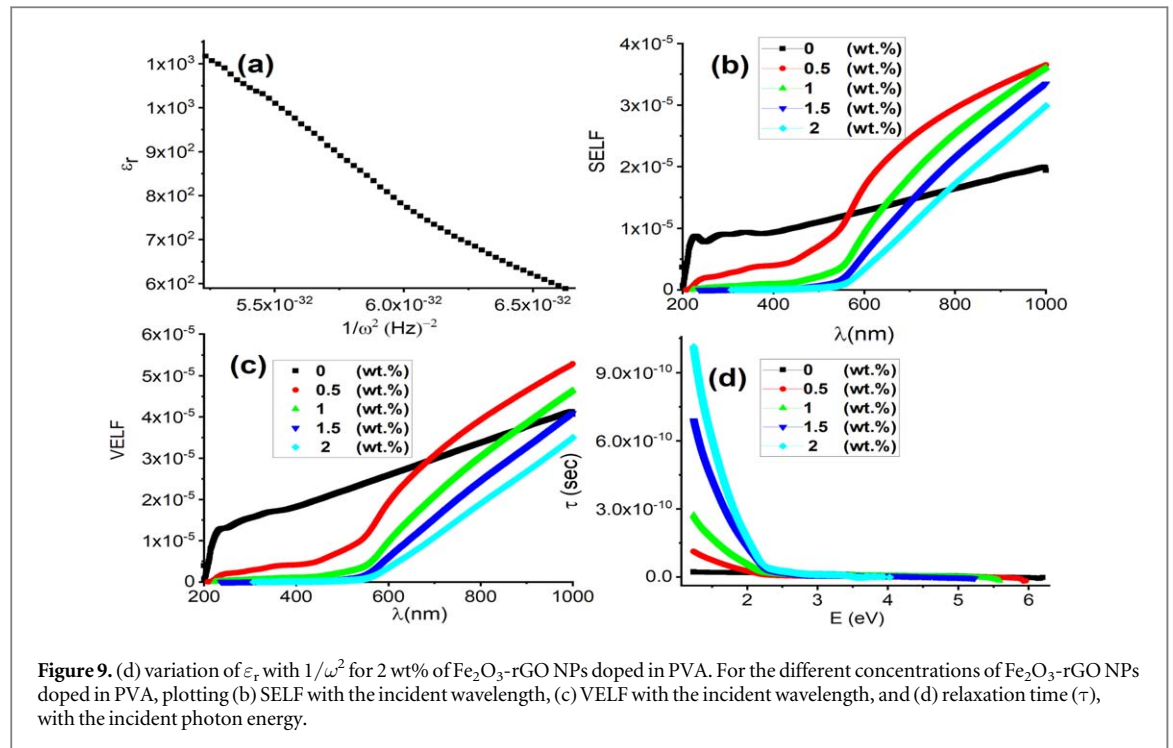
### 3.3.6. High-frequency dielectric constant ( $\epsilon_\infty$ ), and the ratio of ( $N_{\text{opt}}/m^*$ )

The high-frequency dielectric constant ( $\epsilon_\infty$ ) and charge carrier concentration to effective mass ratio ( $N_{\text{opt}}/m^*$ ) were determined by factoring in the influence of free carriers and dispersion vibration modes using the following relation [71]:

$$n^2 = \epsilon_\infty - \frac{1}{4\pi^2\epsilon_0} \left( \frac{e^2}{c^2} \right) \left( \frac{N_{\text{opt}}}{m^*} \right) \lambda^2 \quad (19)$$

Where  $\epsilon_0$  is the free space permittivity,  $N_{\text{opt}}$  is the number of the carrier concentrations,  $m^*$  is the effective mass of the charge carriers, and  $e$  is the electronic charge. Plotting the square of the refractive index ( $n^2$ ) versus the square of wavelength ( $\lambda^2$ ) for 1 wt% of Fe<sub>2</sub>O<sub>3</sub>-rGO NPs doped in PVA, as an example for the others, is depicted in figure 8(d). The values of  $\epsilon_\infty$  deduced from extrapolation of the linear portion of this plot, and the ratio of ( $N_{\text{opt}}/m^*$ ) obtained from the slope were listed in table 3. Shen *et al* [41] determined the effective mass of the charge carriers ( $m^*$ , which equals the production of 0.44 times  $m_0$ ), which was used to calculate the exact values of the charge carrier concentrations ( $N_{\text{opt}}$ ), for all the examined samples. Table 3 shows that increasing the content of





**Figure 9.** (d) variation of  $\epsilon_r$  with  $1/\omega^2$  for 2 wt% of Fe<sub>2</sub>O<sub>3</sub>-rGO NPs doped in PVA. For the different concentrations of Fe<sub>2</sub>O<sub>3</sub>-rGO NPs doped in PVA, plotting (b) SELF with the incident wavelength, (c) VELF with the incident wavelength, and (d) relaxation time ( $\tau$ ), with the incident photon energy.

**Table 3.** Ratio of ( $N_{opt}/m^*$ ), high frequency dielectric constant ( $\epsilon_\infty$ ), optical charge carriers ( $N_{opt}$ ), and plasma frequency ( $\omega_p$ ), for the different nanocomposites.

Fe <sub>2</sub> O <sub>3</sub> - rGO NPs doped in PVA (wt%)	( $N_{opt}/m^*$ ) $\times 10^{60}$ (kg.m <sup>3</sup> ) <sup>-1</sup>	$\epsilon_\infty$ from equation (15)	$N_{opt} \times 10^{30}$ (m <sup>-3</sup> )	$\epsilon_\infty$ from equation (16)	$\omega_p \times 10^{16}$ Hz
0	0.005	5.096	0.002	5.339	0.91
0.5	0.314	95.17	0.126	87.65	2.79
1	1.719	444.1	0.689	438.4	6.98
1.5	6.878	1597	2.757	1550	13.8
2	14.35	3282	5.755	3201	20.0

the doped Fe<sub>2</sub>O<sub>3</sub> - rGO NPs in the host matrix resulted in an improvement in the values of  $\epsilon_\infty$ ,  $N_{opt}$ , and the ratio of ( $N_{opt}/m^*$ ) for all nanocomposites.

### 3.3.7. Plasma frequency and high-frequency dielectric constant ( $\epsilon_\infty$ )

It's worth noting that at the transparency range when the electron damping parameter ( $\gamma$ ) is less than  $\omega$ , the plasma frequency ( $\omega_p$ ), the real part of the dielectric constant ( $\epsilon_r$ ), and the high-frequency dielectric constant ( $\epsilon_\infty$ ) are all related to each other as following [72]:

$$\epsilon_r = \epsilon_\infty - \frac{\omega_p^2}{\omega^2}, \quad \omega_p^2 = \frac{Ne^2}{\epsilon_0 m^*} \quad (20)$$

Figure 9(a) shows the variation of the real part of the dielectric constant ( $\epsilon_r$ ) with  $\omega^{-2}$ , for 2 wt% of Fe<sub>2</sub>O<sub>3</sub>-rGO NPs doped in PVA, as an example for the other studied samples. The values of  $\omega_p^2$ , and  $\epsilon_\infty$  has given from the slope of the straight part of the plot in Figure 9(a), and the intercept, respectively. Table 3 shows the dependence of  $\omega_p$ , and  $\epsilon_\infty$  on the doped amount of Fe<sub>2</sub>O<sub>3</sub>-rGO NPs in PVA. The values of  $\epsilon_\infty$  deduced from equation (15) (5.096, 95.17, 444.1, 1597, and 3282) and those calculated from equation (16) (5.339, 87.65, 438.4, 1550, and 3201) corresponding to 0 wt%, 0.5 wt%, 1 wt%, 1.5 wt%, and 2 wt% of Fe<sub>2</sub>O<sub>3</sub>-rGO NPs doped in PVA, are remarkably close.

### 3.3.8. Energy loss function

It is interesting to know that electrons lose energy as they speed through a material because plasma oscillations in the conduction sea are excited [73]. The functions of volume energy loss (VELF) and surface energy loss (SELF) define the probability that fast electrons will undergo energy loss during their travels within and on the surface of the material, respectively [74]. Real ( $\epsilon_r$ ) and the imaginary ( $\epsilon_i$ ) parts of the dielectric constant were utilized to

calculate the values of (*VELF*) and (*SELF*) as the following [75]:

$$SELF = \frac{\varepsilon_i}{(\varepsilon_r + 1)^2 + \varepsilon_i^2} \quad (21)$$

$$VELF = \frac{\varepsilon_i}{\varepsilon_r^2 + \varepsilon_i^2} \quad (22)$$

For all the studied nanocomposites, the dependence of *SELF* and *VELF* on incident wavelength are shown in Figures 9(b) and (c), respectively. Increasing the content of the doped nanofillers in PVA resulted in lowering the surface and volume energy loss.

### 3.3.9. Dielectric relaxation time

The duration of time that it takes for the electric charges in semiconducting materials to be neutralized during conduction is measured by the dielectric relaxation time ( $\tau$ ). The dielectric relaxation time ( $\tau$ ) is extraordinarily small for metals but is regarded as large for insulators and small for semiconductors. The dielectric relaxation time ( $\tau$ ) was determined according to the following relation [76]:

$$\tau = (\varepsilon_\infty - \varepsilon_r) / \omega \varepsilon_i \quad (23)$$

Figure 9(d) reveals the variation of the dielectric relaxation time ( $\tau$ ) with the incident photon energy for all samples under investigation. Low energy region ( $E < 2.4$  eV) showed dielectric relaxation time- energy dependent behavior, whereas the high energy region ( $E > 2.4$  eV) displayed dielectric relaxation time- energy independent behavior. Also, the dielectric relaxation time ( $\tau$ ) mainly relies on the content of the doped  $\text{Fe}_2\text{O}_3$ -rGO NPs in PVA.

### 3.3.10. Nonlinear optical parameters

Miller's empirical rule was used to calculate the first-ordered nonlinear susceptibility ( $\chi^{(1)}$ ) and the third-ordered nonlinear susceptibility ( $\chi^{(3)}$ ) by using the calculated values of the average refractive index ( $n_{\text{average}}$ ), as follows [77]:

$$\chi^{(1)} = \frac{(-1 + n_{\text{average}}^2)}{4\pi} \quad (24)$$

$$\chi^{(3)} = 1.7 \times 10^{-10} [\chi^{(1)}]^4 \quad (25)$$

Figure 10(a) illustrates the plotting of  $\chi^{(1)}$  with the different contents of  $\text{Fe}_2\text{O}_3$ -rGO NPs doped in PVA. The best fit for the correlation between  $\chi^{(1)}$  and the different contents of  $\text{Fe}_2\text{O}_3$ -rGO NPs is given by the following polynomial equation:

$$Y = A + B_1 X + B_2 X^2 \quad (26)$$

Where Y is  $\chi^{(1)}$ , and the amount of  $\text{Fe}_2\text{O}_3$ -rGO NPs doped in PVA is given by X, A is 0.29,  $B_1$  is 0.05, and  $B_2$  is 0.01.

Figure 10(b) displays the variation of  $\chi^{(3)}$  with the amount of  $\text{Fe}_2\text{O}_3$ -rGO NPs doped in PVA. The following polynomial equation gives the best fit for the relationship between  $\chi^{(3)}$ , and the amount of  $\text{Fe}_2\text{O}_3$ -rGO NPs doped in PVA:

$$Y = A + B_1 X + B_2 X^2 \quad (27)$$

Where Y defines  $\chi^{(3)}$ , and the amount of  $\text{Fe}_2\text{O}_3$ -rGO NPs doped in PVA is given by X, A is  $1.4 \times 10^{-12}$ ,  $B_1$  is  $2.6 \times 10^{-13}$ , and  $B_2$  is  $1.26 \times 10^{-12}$ .

The nonlinear refractive index ( $n_2$ ) was determined by using the following expression for all the examined samples [77]:

$$n_2 = \frac{12\pi\chi^{(3)}}{n_o} \quad (28)$$

Figure 10(c) depicts how the amount of  $\text{Fe}_2\text{O}_3$ -rGO NPs doped in PVA affects the nonlinear refractive index ( $n_2$ ) of the various samples. The following polynomial equation provides the best description of the data in figure 10(c):

$$Y = A + B_1 X + B_2 X^2 \quad (29)$$

Where X describes the amount of  $\text{Fe}_2\text{O}_3$ -rGO NPs doped in PVA, and  $n_2$  is given by Y, A is  $2.4 \times 10^{-11}$ ,  $B_1$  is  $6.6 \times 10^{-12}$ , and  $B_2$  is  $1.6 \times 10^{-11}$ . The values of  $\chi^{(1)}$ ,  $\chi^{(3)}$ , and  $n_2$  were increased by increasing the doped amount of  $\text{Fe}_2\text{O}_3$ -rGO NPs in PVA.

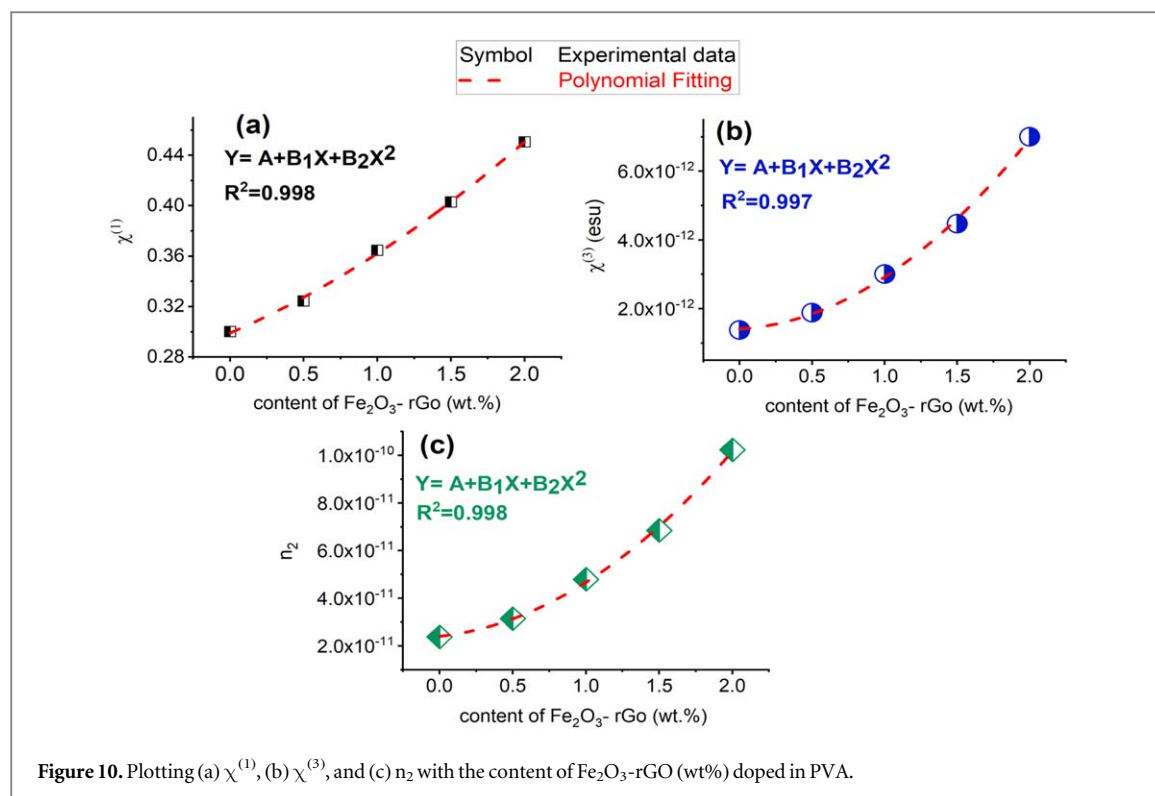


Figure 10. Plotting (a)  $\chi^{(1)}$ , (b)  $\chi^{(3)}$ , and (c)  $n_2$  with the content of Fe<sub>2</sub>O<sub>3</sub>-rGO (wt%) doped in PVA.

#### 4. Conclusion

Fe<sub>2</sub>O<sub>3</sub>-rGO-PVA films have been prepared via the casting method. The structure and optical properties of the PVA matrix have been affected by adding Fe<sub>2</sub>O<sub>3</sub>-rGO NPs. The degree of crystallinity of the polymer films was decreased with the addition of the Fe<sub>2</sub>O<sub>3</sub>-rGO NPs. The morphological analysis shows agglomerations in the PVA matrix with increasing Fe<sub>2</sub>O<sub>3</sub>-rGO NPs concentrations. The transparency of the polymer film decreases with the increase in Fe<sub>2</sub>O<sub>3</sub>-rGO NPs concentration. In addition, the Fe<sub>2</sub>O<sub>3</sub>-rGO NPs additive to the polymer matrix causes the film's transparency in the UV region to vanish, which makes the present polymer films a potential candidate for the UV shielding application. The skin depth of the nanocomposites was decreased with increasing Fe<sub>2</sub>O<sub>3</sub>-rGO NPs in the host matrix due to increasing the absorption and reducing the transparency. The additives of Fe<sub>2</sub>O<sub>3</sub>-rGO NPs to the PVA matrix can increase the conjugation in the monomer unit of the host matrix. The extinction coefficient was found to increase with increasing the wavelength of the incident photons (in the visible region), whereas this behavior disappeared upon introducing Fe<sub>2</sub>O<sub>3</sub>-rGO NPs to the matrix. The optical band gap was found to decrease from 5.956 eV for pure PVA film to 3.076 eV for 2.0 wt% Fe<sub>2</sub>O<sub>3</sub>-rGO NPs in PVA film. Moreover, the band tail width was found to increase with the increase in Fe<sub>2</sub>O<sub>3</sub>-rGO NPs in the PVA matrix, indicating the rise in polymer disordering and the decrease in the optical band gap. The values of  $\chi^{(1)}$ ,  $\chi^{(3)}$ , and  $n_2$  were increased by increasing the doped amount of Fe<sub>2</sub>O<sub>3</sub>-rGO NPs in PVA.

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#### Data availability statement

All data that support the findings of this study are included within the article. The data that support the findings of this study are available upon reasonable request from the authors.

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